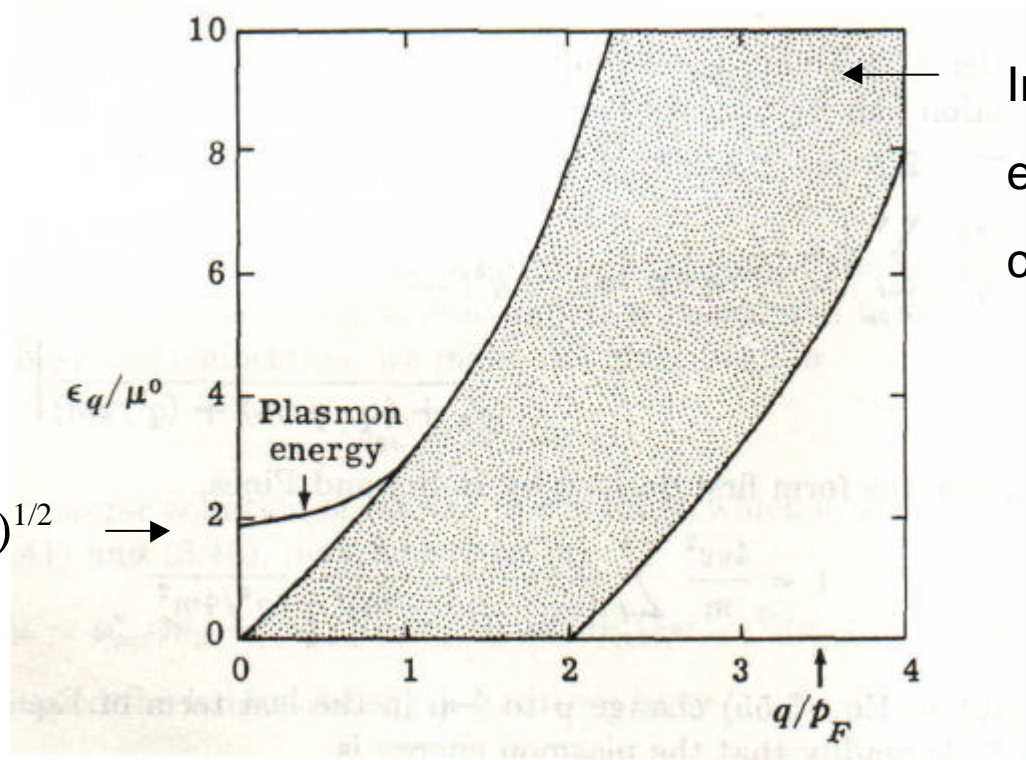


Collective excitations  
in  
integer quantum Hall system

師範大學物理系  
張明哲

Plasma excitation in **3DEG**, NO magnetic field

$r_s=4$ , sodium



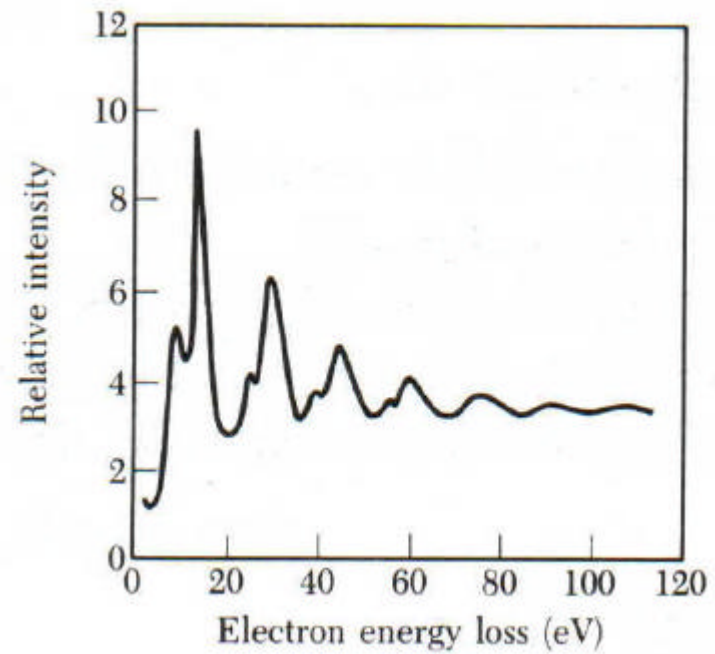
Independent  
electron-hole  
continuum

$$\omega_p = (4\pi\rho e^2/m)^{1/2}$$

**2DEG**, NO magnetic field ? gapless excitation

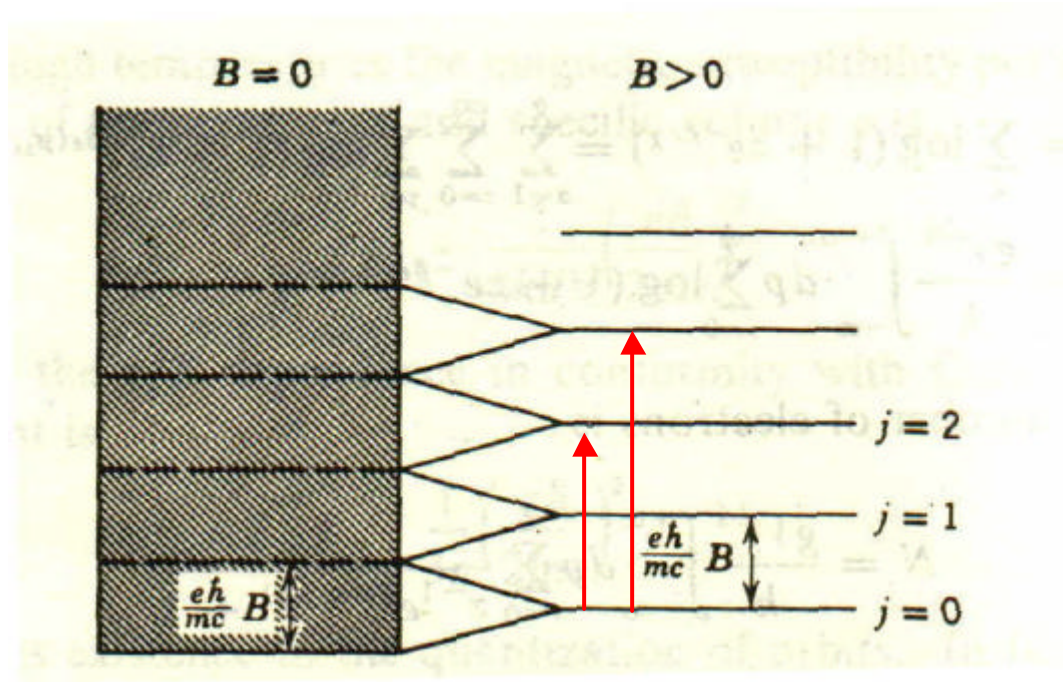
## Experimental observation

Al: 10.3 eV + 15.3 eV



10.3 eV due to surface plasmons

2DEG in a strong magnetic field  
emergence of Landau levels

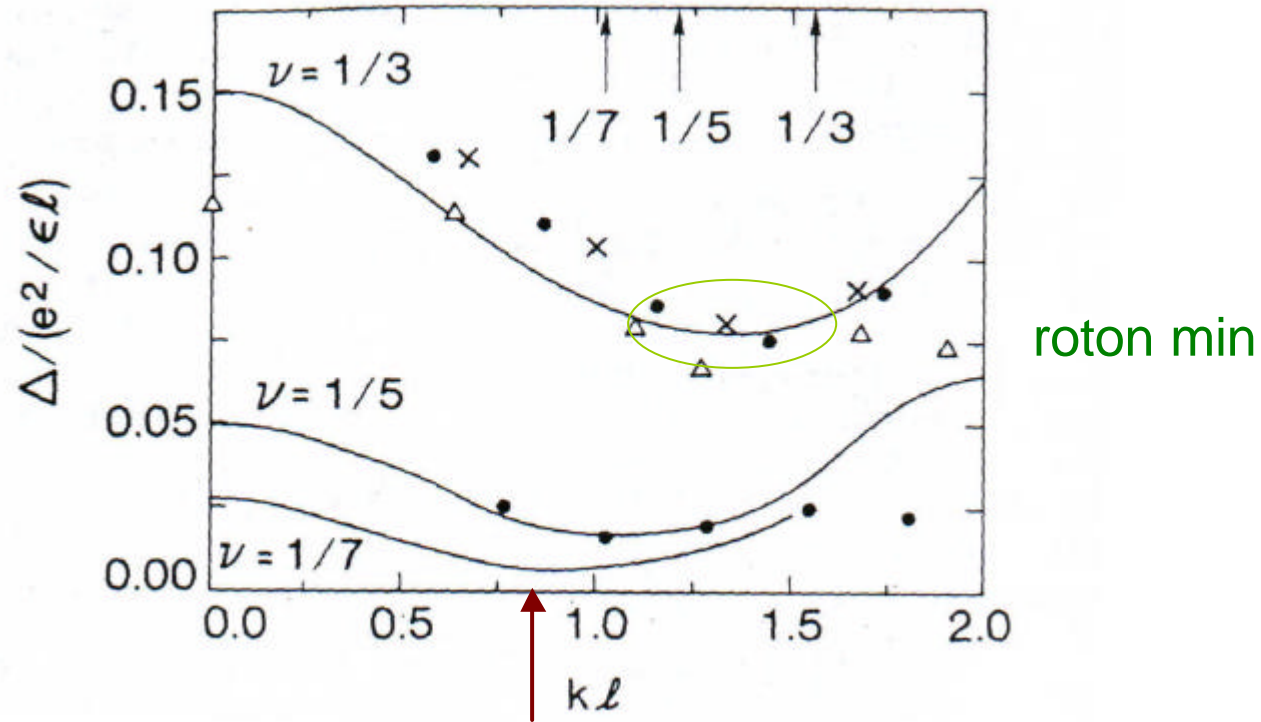


- magnetoplasma modes: inter-LL excitations
- can be Integer filling / non-integer filling
- Inter-LL excitation/intra-LL excitation
- Charge excitation / spin excitation

Fractional case, intra-LL charge excitations

Finite energy gap stabilizes the FQHL

Magneton mode similar to rotons in superfluid



indication of Wigner crystal instability?

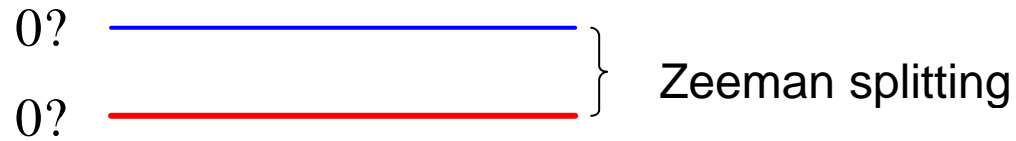
(Girvin/MacDonald/Platzman 1985 PRL)

# Integer case

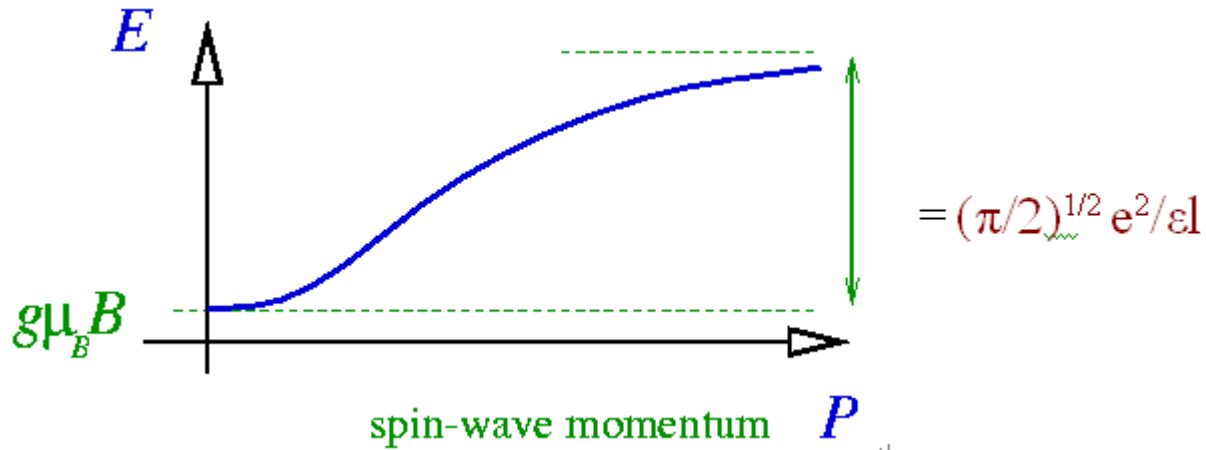
classification of collective excitations

Intra-LL :

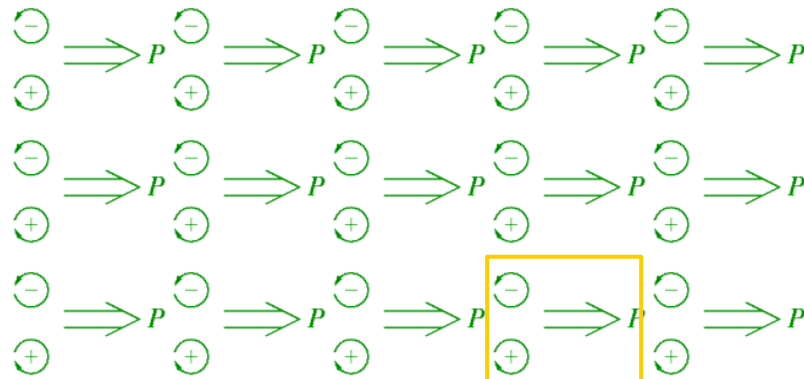
Spin flip excitations only



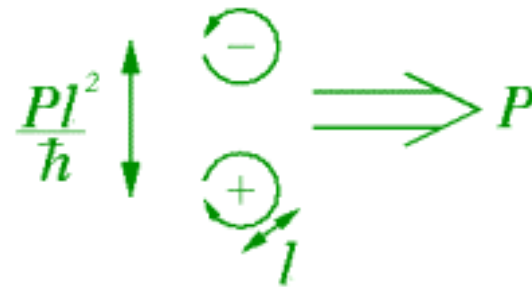
energy dispersion of spin wave



spin wave as coherent superposition of e-h pairs



semiclassical picture



$$E_{\infty} = (\pi/2)^{1/2} e^2/\epsilon l$$

= the energy to flip the spin of an electron

and take it far away from the hole

= the exchange energy of an electron in the LLL

## Excitation energy

Single particle:    cyclotron energy  
                          Zeeman energy

Manybody effect: exchange interaction  
                          polarization and screening  
                          excitonic attraction ... etc

## Methods of calculation

Equation of motion method

→ Diagrammatic expansion

    The pole of response function

    → energy spectrum of collective excitation

    Time dependent HFA

# Calculation of response function

$$\chi(\vec{q}, \omega) = \text{---} \rightarrow \text{---} \text{---} \text{---} \rightarrow \text{---} \text{---} \text{---} \text{---} \rightarrow$$

## Renormalized electron propagator

$$G = \text{---} \rightarrow \text{---} \cong \text{---} \rightarrow \text{---} + \text{---} \rightarrow \text{---} \text{---} \text{---} \rightarrow$$

## Bethe-Salpeter equation

$$\Gamma = \text{---} \text{---} \text{---} \text{---} \rightarrow \text{---} \text{---} \text{---} \text{---} \rightarrow \cong \text{---} \text{---} \text{---} \text{---} \rightarrow \text{---} \text{---} \text{---} \text{---} \rightarrow + \text{---} \text{---} \text{---} \text{---} \text{---} \text{---} \text{---} \text{---} \rightarrow \text{---} \text{---} \text{---} \text{---} \rightarrow$$

ladder

bubble

## Response function

$$\chi(k, i\omega) \equiv - \int_0^\infty d\tau e^{i\omega\tau} \langle T_\tau \hat{\rho}(\vec{k}, \tau) \hat{\rho}(-\vec{k}, 0) \rangle.$$

$$= \underbrace{\sum_{n_\alpha n_\beta} \sum_{m_\alpha m_\beta} \langle n_\alpha, m_\alpha | e^{i\vec{k} \cdot \vec{r}} | n_\beta, m_\beta \rangle}_{\text{inter/intra-LL indices}} \underbrace{D_{n_\alpha n_\beta}(i\omega)}_{\text{free 2-particle propagator}} \Gamma_{n_\alpha n_\beta}^{m_\alpha m_\beta}(k, i\omega)$$

free 2-particle propagator

$$\underline{D_{n_\alpha' n_\beta'}(i\omega)} \equiv \int \frac{d\omega'}{2\pi} G_{n_\alpha'}(i\omega' - i\omega) G_{n_\beta'}(i\omega')$$

## BS equation for the vertex

$$\Gamma_{n_\alpha n_\beta}(k, i\omega) = \frac{i^{n_\alpha n_\beta}}{2\pi} e^{-k^2/2} \bar{g}_{n_\alpha n_\beta}(-\vec{k})$$

$$- \sum_{n_\alpha' n_\beta'} \left[ \underbrace{V_{n_\alpha n_\beta}^{n_\alpha' n_\beta'}(k)}_{\text{ladder}} - \underbrace{U_{n_\alpha n_\beta}^{n_\alpha' n_\beta'}(k)}_{\text{bubble}} \right] \underline{D_{n_\alpha' n_\beta'}(i\omega)} \Gamma_{n_\alpha' n_\beta'}(k, i\omega)$$

matrix form of the BS equation

$$\sum_{n_\alpha, n_\beta} \left[ \delta_{n_\alpha, n_\alpha'} \delta_{n_\beta, n_\beta'} D_{n_\alpha', n_\beta'}^{-1}(i\omega) + V_{n_\alpha n_\beta}^{n_\alpha' n_\beta'}(k) - U_{n_\alpha n_\beta}^{n_\alpha' n_\beta'}(k) \right] D_{n_\alpha', n_\beta'}(i\omega) \Gamma_{n_\alpha', n_\beta'}(k, i\omega) = \frac{i^{n_\alpha \beta}}{2\pi} e^{-k^2/2} \bar{g}_{n_\alpha n_\beta}(-\bar{k}),$$

$$D_{n_\alpha', n_\beta'}^{-1}(i\omega) = \left( i\omega - (n_\alpha' - n_\beta')\omega_C - \left( \Sigma_{n_\alpha'}^{n_0}, -\Sigma_{n_\beta'}^{n_0} \right) \right)$$

↑  
self-energy

bubble (Coulomb local field)

$$U_{n_\alpha n_\beta}^{n_\alpha' n_\beta'}(k) \equiv i^{n_\alpha \beta - n_\alpha \beta'} e^{-k^2/2} \bar{g}_{n_\alpha n_\beta}(-\bar{k}) \frac{\tilde{v}(k)}{2\pi} g_{n_\alpha' n_\beta'}(-\bar{k})$$

ladder (exchange local field)

$$V_{n_\alpha n_\beta}^{n_\alpha' n_\beta'}(k) \equiv 2\pi \int d^2 r_1 \int d^2 r_2 \Phi_{n_\alpha, n_\beta}^\kappa(\vec{r}_1, \vec{r}_2) v(\vec{r}_1 - \vec{r}_2) \bar{\Phi}_{n_\alpha', n_\beta'}^\kappa(\vec{r}_1, \vec{r}_2)$$

exciton wave function ↑

poles of the response function  
= eigenvalues of the matrix eq.

consider ? n=1, neglect inter-mode coupling  
one of the diagonal element

$$M_{2,1}^{2,1}(k, \omega) = \omega - \omega_C - \frac{e^2}{\epsilon l_0} (\tilde{\Sigma}_{n=2}^1 - \tilde{\Sigma}_{n=1}^1) - \frac{e^2}{\epsilon l_0} \tilde{U}_{2,1}^{2,1}(k) + \frac{e^2}{\epsilon l_0} \tilde{V}_{2,1}^{2,1}(k) = 0.$$

transition energy =

cyclotron energy + self energy correction (k independent)

+ Coulomb local field correction

+ exchange local field correction

## General behavior of the manybody corrections

$k \rightarrow 0$

- { bubble correction  $\rightarrow 0$
- { ladder correction  $\rightarrow$  self-energy correction

Kohn theorem: the cyclotron energy gap is NOT

renormalized by many-body correction

$k \rightarrow \infty$

- { bubble correction  $\rightarrow 0$  exponentially
- { ladder correction  $\rightarrow 0$  algebraically

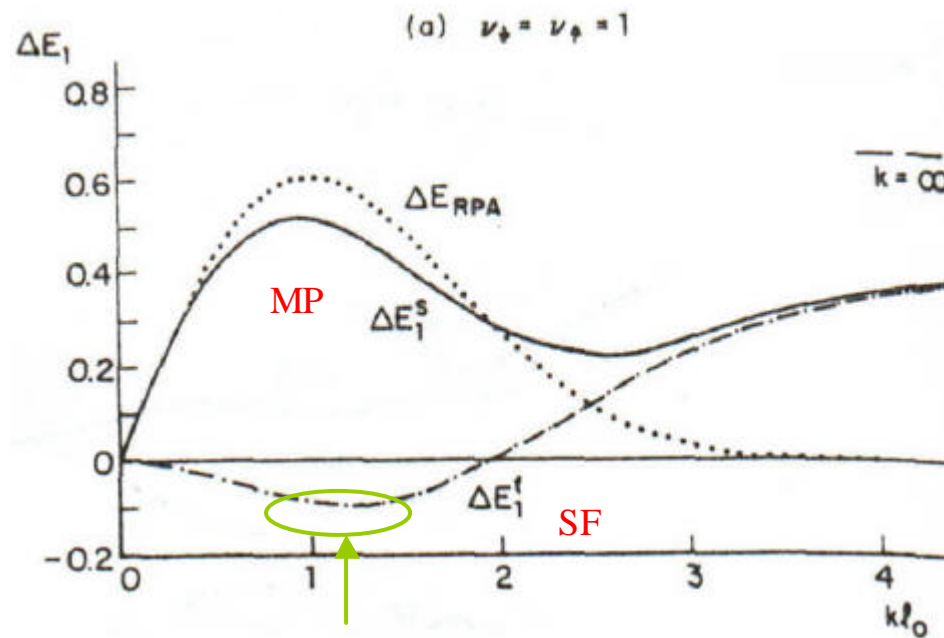
Coulomb interaction prohibits spin-flip

$\rightarrow$  no bubble correction for spin-flip excitations

energy dispersion determined solely by

ladder correction

$$\nu = 2$$



roton feature (spin triplet roton)