

# Solid state physics 簡答

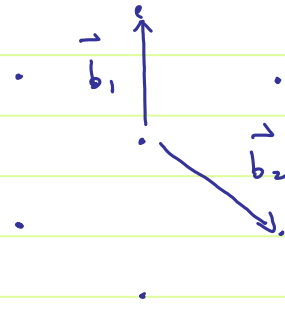
1. (a)

$$\begin{cases} \vec{a}_1 = a \hat{x} \\ \vec{a}_2 = \frac{a}{2} \hat{x} + \frac{\sqrt{3}}{2} a \hat{y} \end{cases}$$

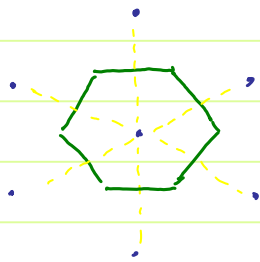
$$\vec{b}_1 \cdot \vec{a}_1 = 2\pi, \quad \vec{b}_1 \cdot \vec{a}_2 = 0$$

$$\vec{b}_2 \cdot \vec{a}_1 = 0, \quad \vec{b}_2 \cdot \vec{a}_2 = 2\pi$$

$$\rightarrow \begin{cases} \vec{b}_1 = \frac{2\pi}{a} \hat{x} - \frac{2\pi}{\sqrt{3}a} \hat{y} \\ \vec{b}_2 = \frac{4\pi}{\sqrt{3}} a \hat{y} \end{cases}$$

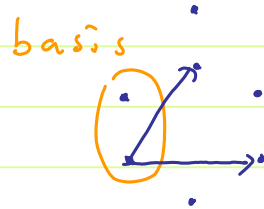


(b)



(c)

$$\begin{cases} \vec{d}_1 = 0 \\ \vec{d}_2 = \frac{a}{\sqrt{3}} \hat{y} \end{cases}$$



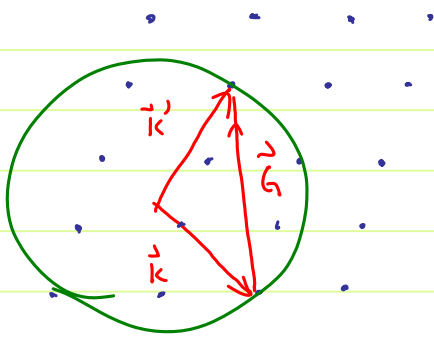
$$\begin{aligned} S(\vec{G}) &= f_a \sum_j e^{-i\vec{G} \cdot \vec{d}_j} \\ &= f_a (1 + e^{-i\vec{G} \cdot \vec{d}_2}) \end{aligned}$$

$$\begin{aligned} \vec{G} \cdot \vec{d}_2 &= (h \vec{b}_1 + k \vec{b}_2) \cdot \vec{d}_2 \\ &= \left[ \frac{2\pi}{a} h \hat{x} + \left( \frac{4\pi}{\sqrt{3}a} k - \frac{2\pi}{\sqrt{3}a} h \right) \hat{y} \right] \cdot \frac{a}{\sqrt{3}} \hat{y} \\ &= \frac{2\pi}{3} (2k - h) \end{aligned}$$

$$\therefore S_{kh} = f_a \left( 1 + e^{-2\pi i \frac{2k-h}{3}} \right)$$

$$|S_{kh}|^2 = 2|f_a|^2 \left[ 1 + \cos \frac{2\pi}{3} (2k-h) \right] = \begin{cases} 4|f_a|^2 & (0,0), (2,1) \\ |f_a|^2 & \dots \\ & (1,0), (0,1) \\ & \dots \end{cases} \quad \text{for } (h,k) =$$

2 (a)

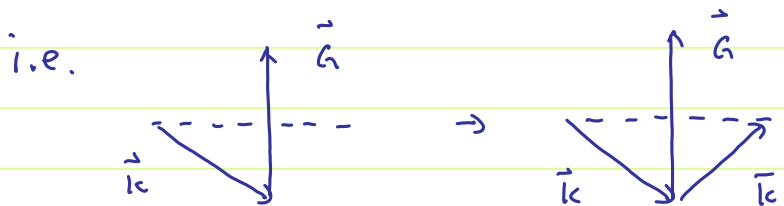


attach  $\vec{k}$  to a lattice point  
draw a circle (see fig.)  
if the circle (or sphere, in 3D)  
intersects w/ a point, then  
we find a  $\vec{k}'$  that satisfies  
the Laue condition.

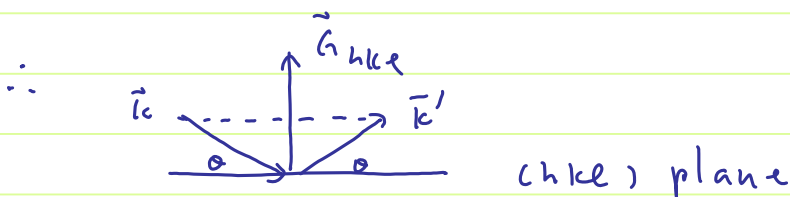
(b) Laue condition:

$$\vec{k}' = \vec{k} + \vec{G}$$

$$\rightarrow \vec{k} \cdot \vec{G} = -\frac{G^2}{2}$$



$\vec{G}_{hkl} \perp (hkl)$  plane,  $|\vec{k}| = |\vec{k}'|$



from the geometry, we get

$$2k \sin \theta = G_{hkl} = \frac{2\pi}{d_{hkl}} n$$

$$k = \frac{2\pi}{\lambda}$$

$$\therefore 2 d_{hkl} \sin \theta = n \lambda$$

3. (a) A fermi disk has radius  $k_F$

$$2 \frac{\pi k_F^2}{(2\pi/L)^2} = N$$

$$\rightarrow \frac{k_F^2}{2\pi} = \frac{N}{L^2} = n$$

$$\rightarrow k_F = \sqrt{2\pi n}$$

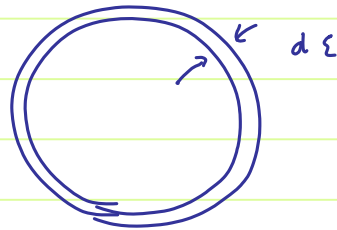
(b)

$$D(\epsilon) d\epsilon = 2 \cdot \frac{2\pi k dk}{(2\pi/L)^2}$$

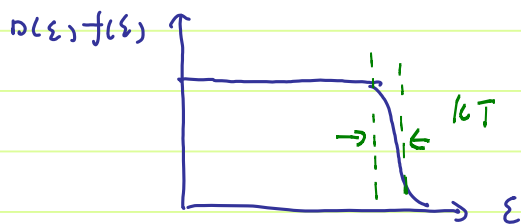
$$= \frac{L^2}{\pi} k dk$$

$$\frac{\hbar^2 k^2}{2m} = \epsilon \rightarrow \frac{\hbar^2}{m} k dk = d\epsilon$$

$$\therefore \underline{D(\epsilon) d\epsilon} = \underline{\frac{L^2}{\pi} \frac{m}{\hbar^2} d\epsilon}$$



(c) in 2-dim



thermally excited electron  $\sim N \frac{kT}{\epsilon_F}$

each electron has energy  $\sim kT$

$\therefore$  electrons absorb thermal energy  $\Delta U \sim N \frac{(kT)^2}{\epsilon_F}$ .

specific heat  $C_e = \frac{dU}{dT} \sim N \frac{kT}{\epsilon_F} \left( \text{or } R \frac{T}{T_F} \right)$

$\propto T^1$  as in 3D

4. (a) # of k points in the 1st BZ

$$= \frac{1 \text{ nm}^3}{5 \text{ \AA}^3} \quad 1 \text{ nm} = 10^7 \text{ \AA}$$

$$= 2 \times 10^{20}$$

each band can host

$$2 \times 2 \times 10^{20} = 4 \times 10^{20} \text{ electrons}$$

(b)

$$\vec{v} = \langle \psi_{nk} | \frac{\vec{p}}{m} | \psi_{nk} \rangle$$

$$= \langle u_{nk} | \frac{\vec{p} + \hbar \vec{k}}{m} | u_{nk} \rangle$$

$$= \langle u_{nk} | \frac{\partial \tilde{H}}{\partial \hbar \vec{k}} | u_{nk} \rangle, \quad \tilde{H} = \frac{1}{2m} (\vec{p} + \hbar \vec{k})^2 + U(\vec{r})$$

$$= \frac{\partial}{\partial \hbar \vec{k}} \langle u_{nk} | \tilde{H} | u_{nk} \rangle - \left( \frac{\partial}{\partial \hbar \vec{k}} u_{nk} | \tilde{H} | u_{nk} \right) - \langle u_{nk} | \tilde{H} | \frac{\partial u_{nk}}{\partial \hbar \vec{k}} \rangle$$

$$= \frac{\partial \epsilon_{nk}}{\partial \hbar \vec{k}} - \epsilon_{nk} \frac{\partial}{\partial \hbar \vec{k}} \langle u_{nk} | u_{nk} \rangle$$

$$= \frac{\partial \epsilon_{nk}}{\partial \hbar \vec{k}}$$